

Magnetic response of the J_1 - J_2 Spin Hamiltonian and its Implication for FeAs

arXiv:0904.3809

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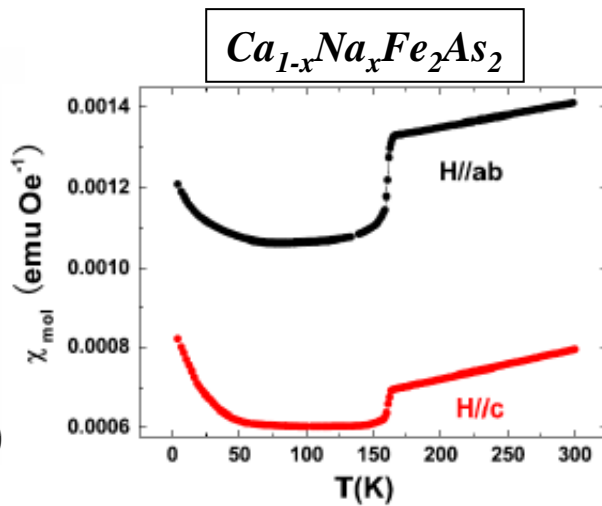
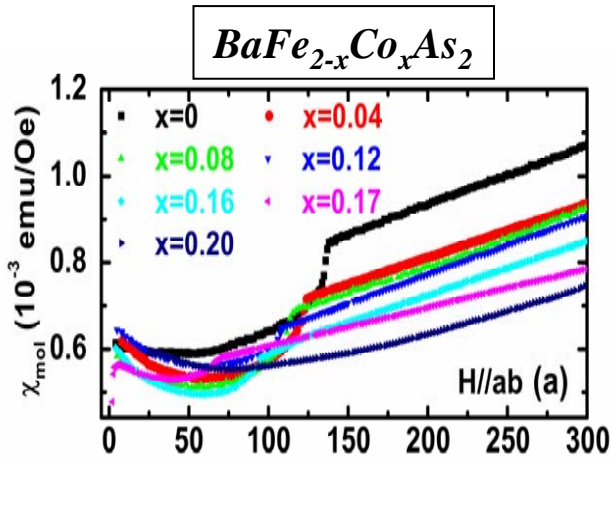
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6/2009 @ Gordon Research Conference, HK

Magnetic properties of parent FeAs compounds

- Uniform magnetic susceptibility above SDW ordering (<200K)
- **Experiments:** T-linear over, *ie.*, 150~500-700K!



X.F.Wang, *et al.* NJP. 11, 0405(09)

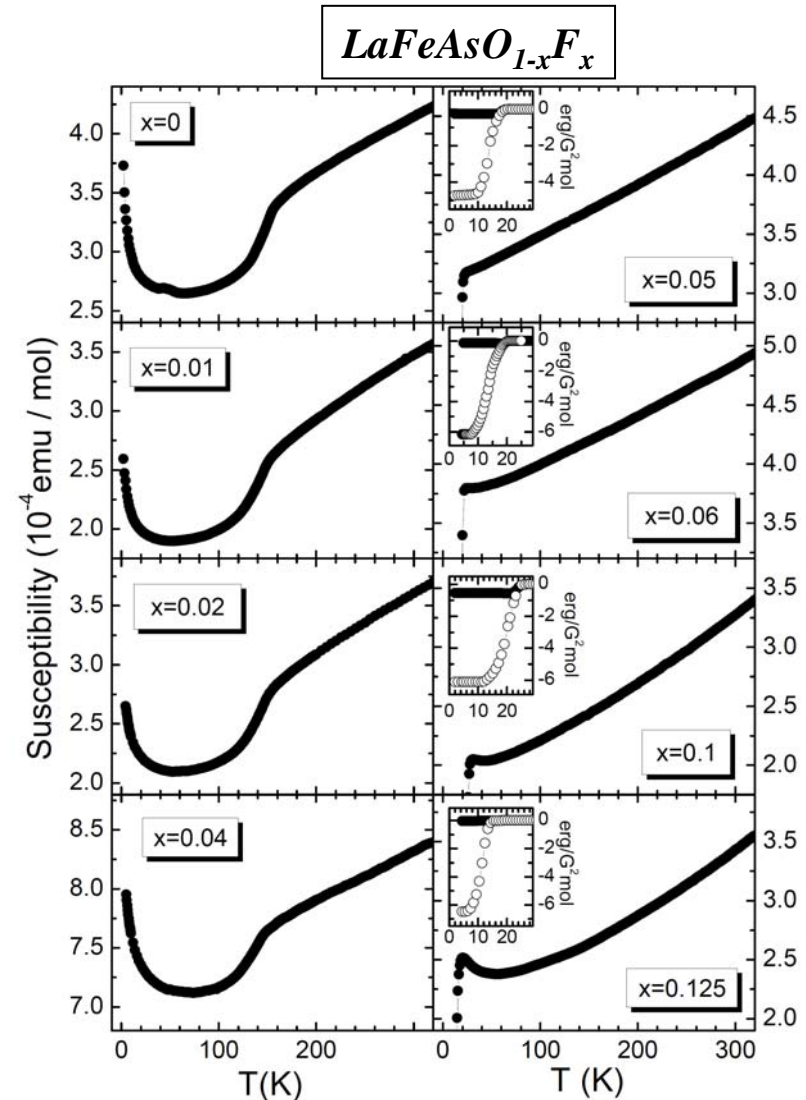
G. Wu, *et al.* JPhysC.20, 422201 (09)

- In disagreement with Pauli-like (itinerant) or Curie-Weiss-like (local)

- $T < T_{SDW}$, AF SDW magnetic order: $(\pi, 0)$

C.de la Cruz, *et al.*, Nature(08)

- $T > T_{SDW}$, no global AF, *a form of "correlated paramagnet"?*



R.Klingeler, *et al.*, arXiv:0808.0708v1

Motivation

• J_1 - J_2 Spin Hamiltonian

$$H \sim \sum_{ij} J_{ij}^{\alpha\beta} S_{i\alpha} \cdot S_{j\beta}$$

(Si & Abrahams, PRL(08))

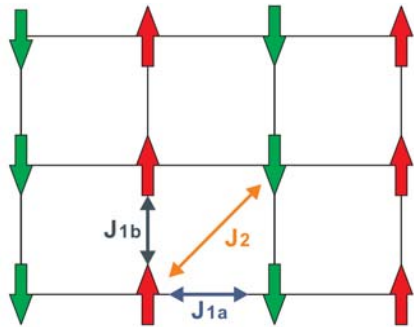
Disregarding

multi-orbital nature

$$H = J_1 \sum_{\langle ij \rangle} S_i \cdot S_j + J_2 \sum_{\langle\langle ij \rangle\rangle} S_i \cdot S_j$$

(Ma, Lu & Xiang, PRB (08))

• Physical $2J_2/J_1$ parameters: first-principles calculation by Han, Yin, Pickett & Savrasov, PRL (09))



Fe moments (in B),
in-plane exchange interactions (in meV)
using experimental $z(\text{As})$.

$$2J_2/J_{1a}: 0.64 \sim 1.05$$

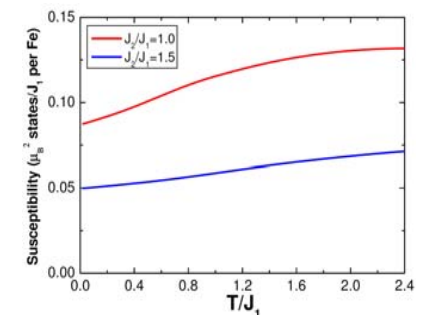
System	Moment	J_{1a}	J_2	J_{1b}	$J_{1a}/2J_2$	$J_{1a} + 2J_2$
LaFeAsO	1.69	47.4	22.4	-6.9	1.06	92.2
CeFeAsO	1.79	31.6	15.4	2.0	1.03	62.4
PrFeAsO	1.76	57.2	18.2	3.4	1.57	93.6
NdFeAsO	1.49	42.1	15.2	-1.7	1.38	72.5
CaFe ₂ As ₂	1.51	36.6	19.4	-2.8	0.95	75.4
SrFe ₂ As ₂	1.69	42.0	16.0	2.6	1.31	74.0
BaFe ₂ As ₂	1.68	43.0	14.3	-3.1	1.51	71.5
KFe ₂ As ₂	1.58	42.5	15.0	-2.9	1.42	72.5
LiFeAs	1.69	43.4	22.9	-2.5	0.95	89.2

• Spin Hamiltonian for $T > T_{\text{SDW}}$, preformed moments exist (G.M.Zhang, *et al.* EPL(09))

- Linear-T uniform susceptibility

ie. Zhang *et al.*(EPL,09) $\chi_u = \chi_0(1 + a(\frac{T}{J_1}))$ for $J_2/J_1=1, 1.5$

Kou *et al.*(arXiv:0811.0411v3) $\chi_u = \chi_{lo} + \chi_{it}$



How about the uniform and staggered susceptibilities in the whole range $2J_2/J_1$?

Abstract

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Classical Monte Carlo method

Schwinger boson mean field theory



Antiferromagnetic J_1 - J_2 Spin Hamiltonian



$$\chi_k^\alpha = \frac{\langle S_{-k}^\alpha S_k^\alpha \rangle - \langle S_{-k}^\alpha \rangle \langle S_k^\alpha \rangle}{TN}$$

k = any vector !!

$$S_k^\alpha = \sum_i S_i^\alpha e^{-ik \cdot r_i}$$

$$\alpha = x, y, z$$

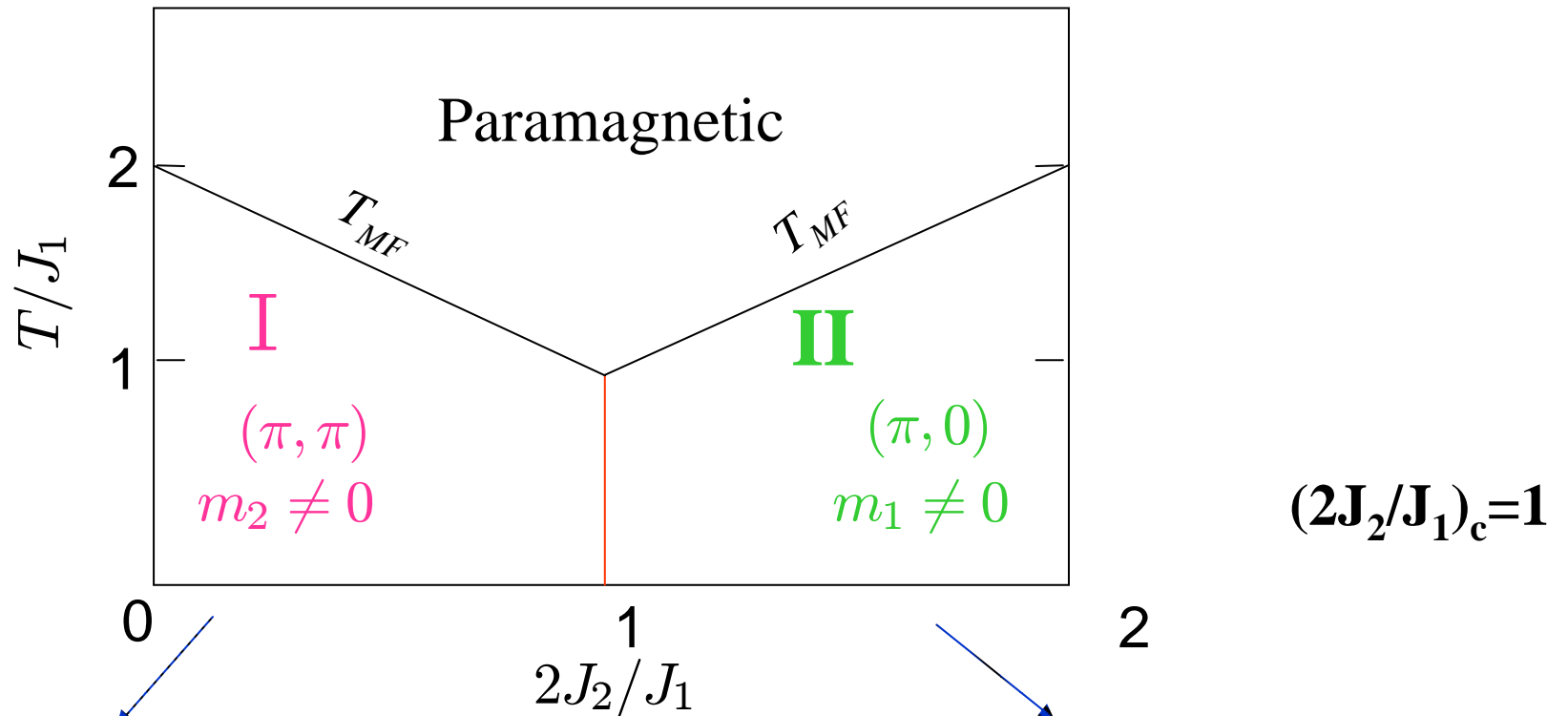
$N = L \times L$ is lattice size

Magnetic susceptibilities at wave vectors : $k = (0, 0), (\pi, 0), (0, \pi), (\pi, \pi)$

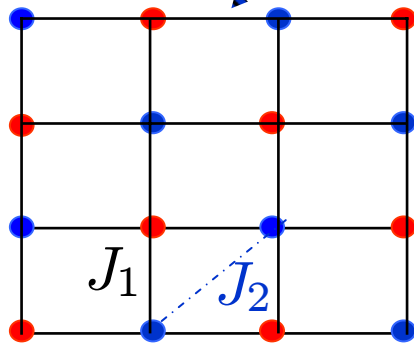
- A nearly linear-T behavior for the uniform susceptibility
- Implication for parent FeAs compounds

Review of Classical Mean Field Phase Diagram

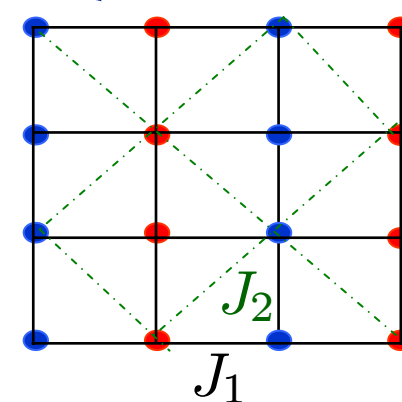
Ordering pattern: $\langle S_i \rangle = m_1(-1)^{x_i} + m_2(-1)^{x_i+y_i}$



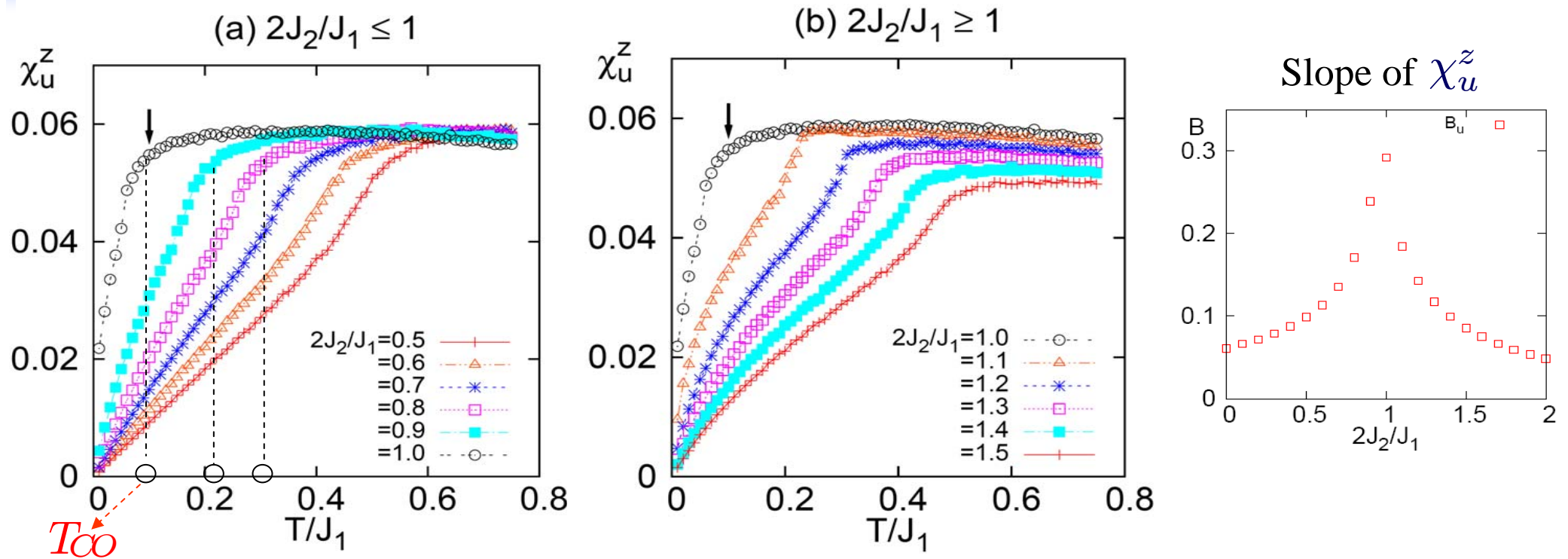
AF1
 (π, π)



AF2
 $(\pi, 0)$



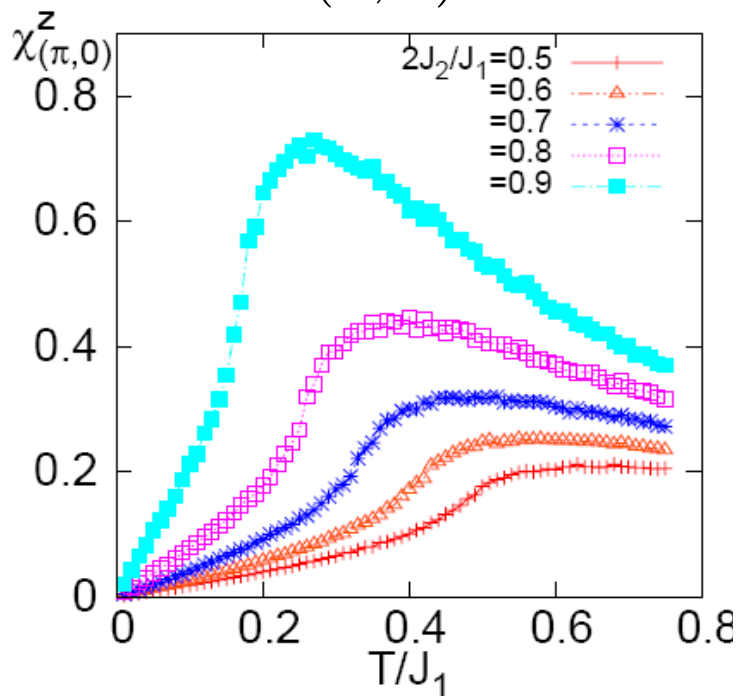
Uniform susceptibility (classical Monte Carlo(CMC))



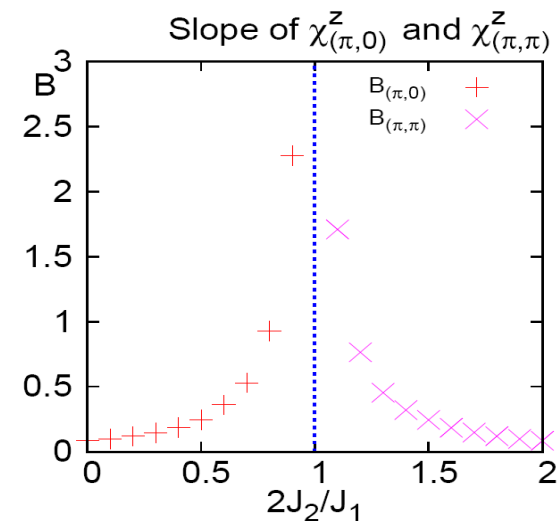
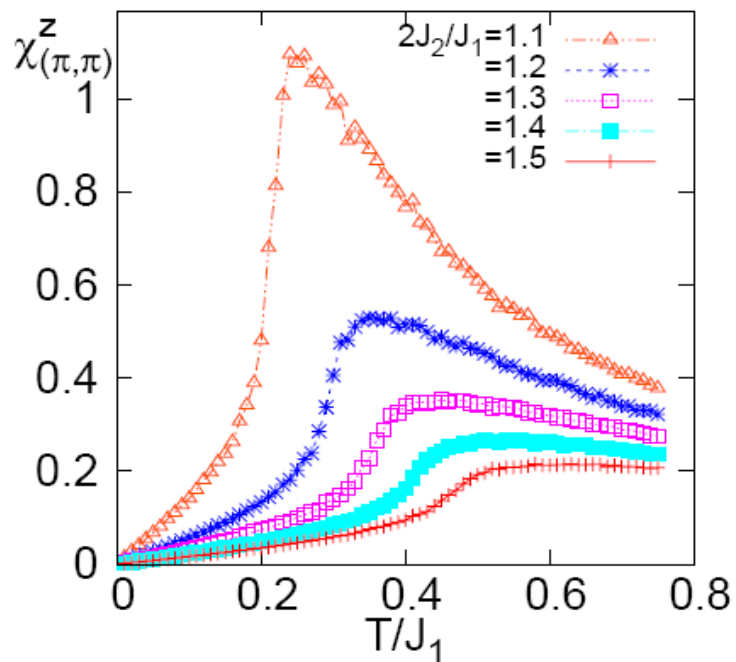
- The linear-T behavior of χ_u^z is a natural aspect of the J_1 - J_2 spin model ;
- The crossover temperature $T_{CO} \sim 0.8J_1 \sim 400K$;
- The linearity of χ_u^z persists over a wider temperature range when the system is away from the critical value.

Staggered susceptibility (CMC)

In (π, π) order



In $(\pi, 0)$ order



- Non-dominant, non-divergent susceptibilities , roughly linear
- A linear-T behavior with slope B

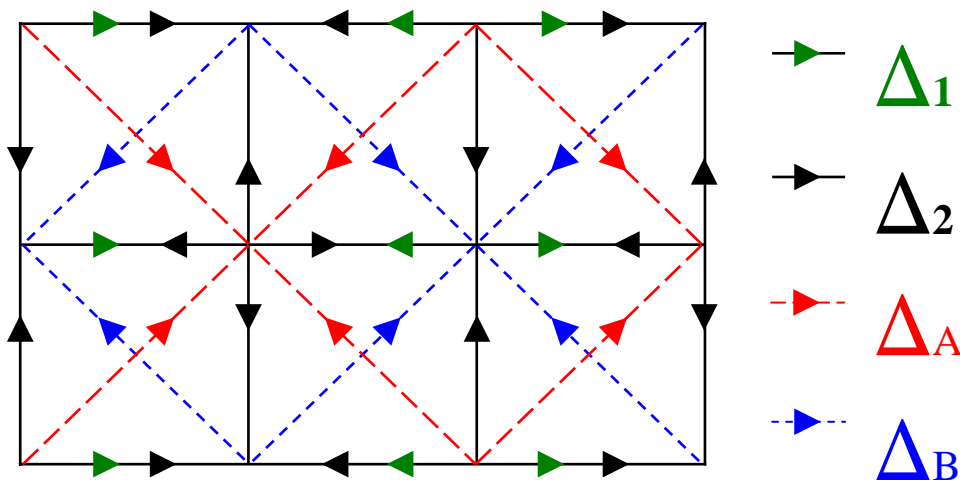
Schwinger Boson mean field theory(SBMFT)

• J_1 - J_2 Hamiltonian in mean field form

$$\vec{S} = \frac{1}{2} b_\alpha^\dagger \vec{\sigma}_{\alpha\beta} b_\beta$$

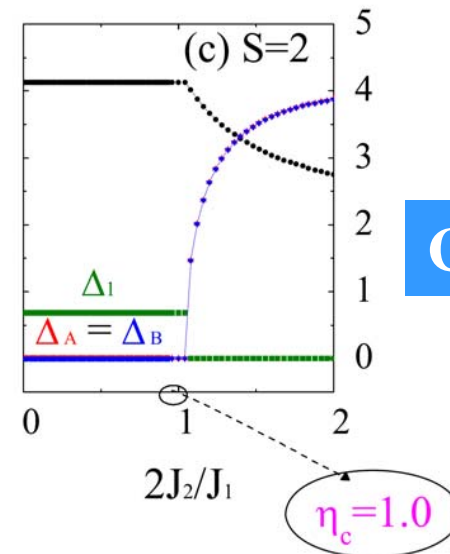
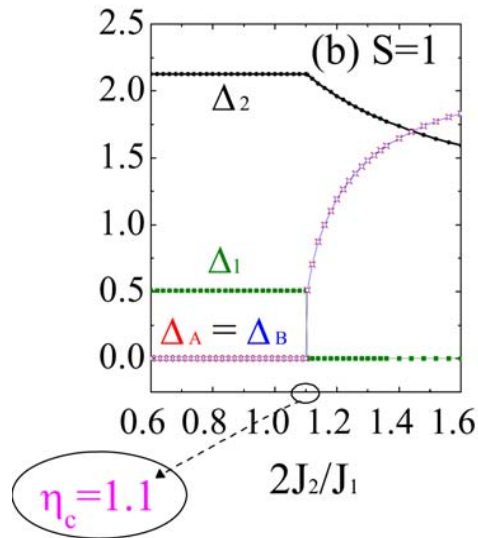
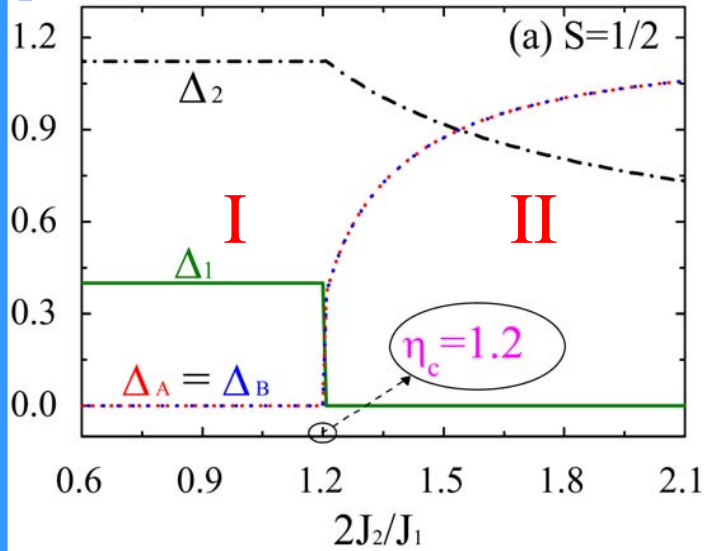
$$H_{MF} = -\frac{J_1}{2} \sum_{\langle ij \rangle} \Delta_{ij}^* A_{ij} - \frac{J_2}{2} \sum_{\langle\langle ij \rangle\rangle} \Delta_{ij}^* A_{ij} + h.c. + \sum_i \lambda_i (b_{i1}^\dagger b_{i1} + b_{i2}^\dagger b_{i2} - 2S)$$

where the operators are $A_{ij} = b_{i1} b_{j2} - b_{i2} b_{j1} = -A_{ji}$, $\Delta_{ij} = \langle A_{ij} \rangle$.



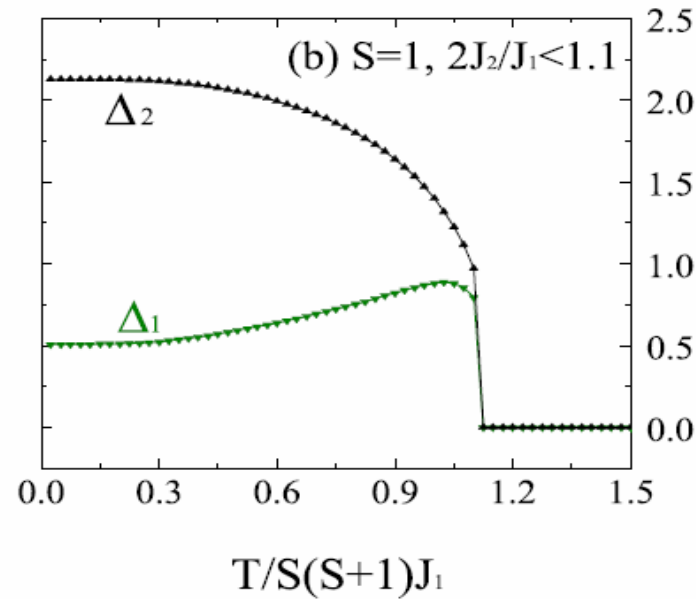
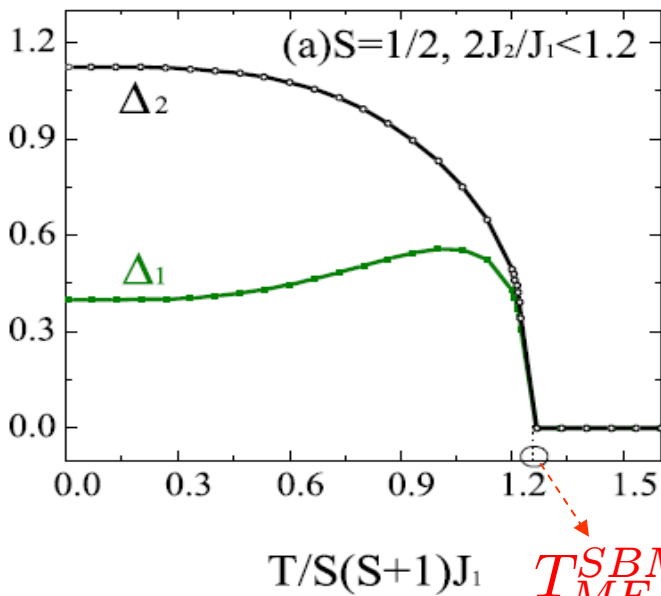
k	Gap parameters
(π, π)	Δ_2
$(\pi, 0)$	$\Delta_{1,A,B}$

Gap parameters(SBMFT)



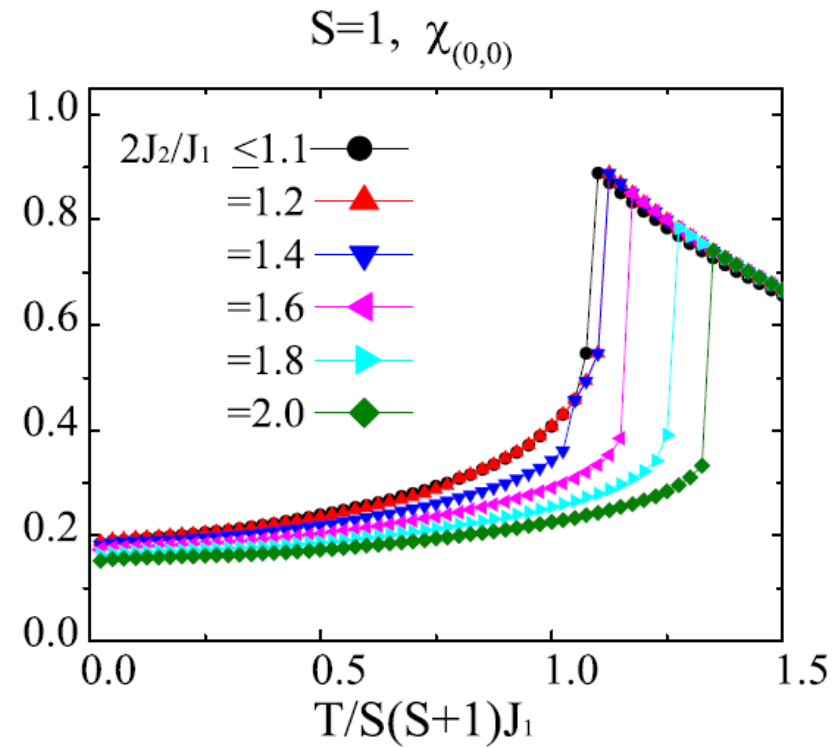
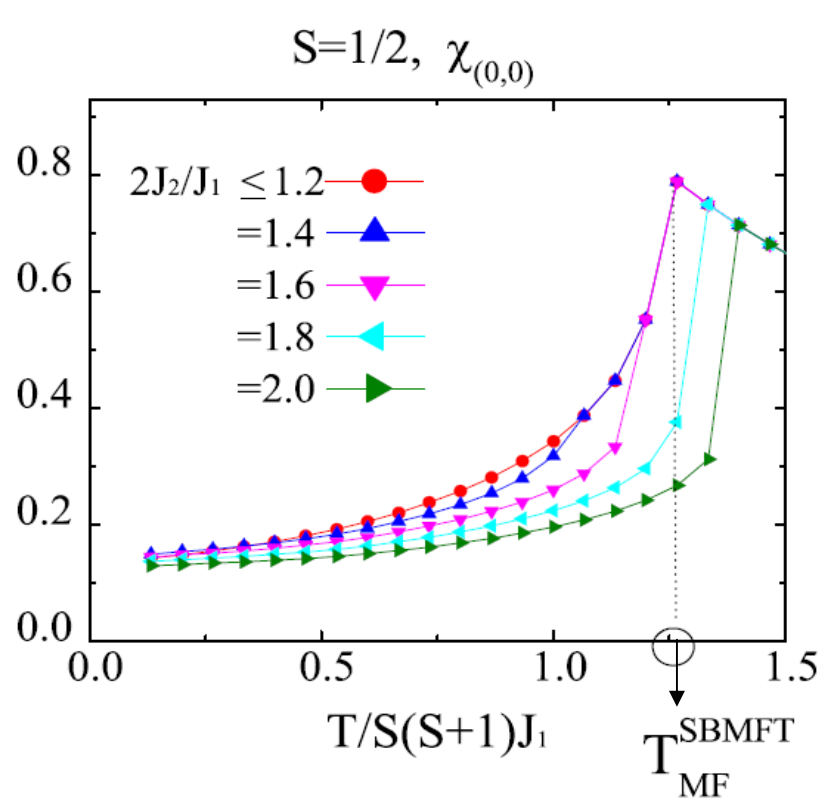
Critical point η_c

consistent with Capriotti, PRL (04)



$T > T_{CO}^{SBMFT}$
Gap parameters=0

Uniform susceptibility(SBMFT)



- **Linear-T behavior,** $\chi_{(0,0)} \sim A + BT$

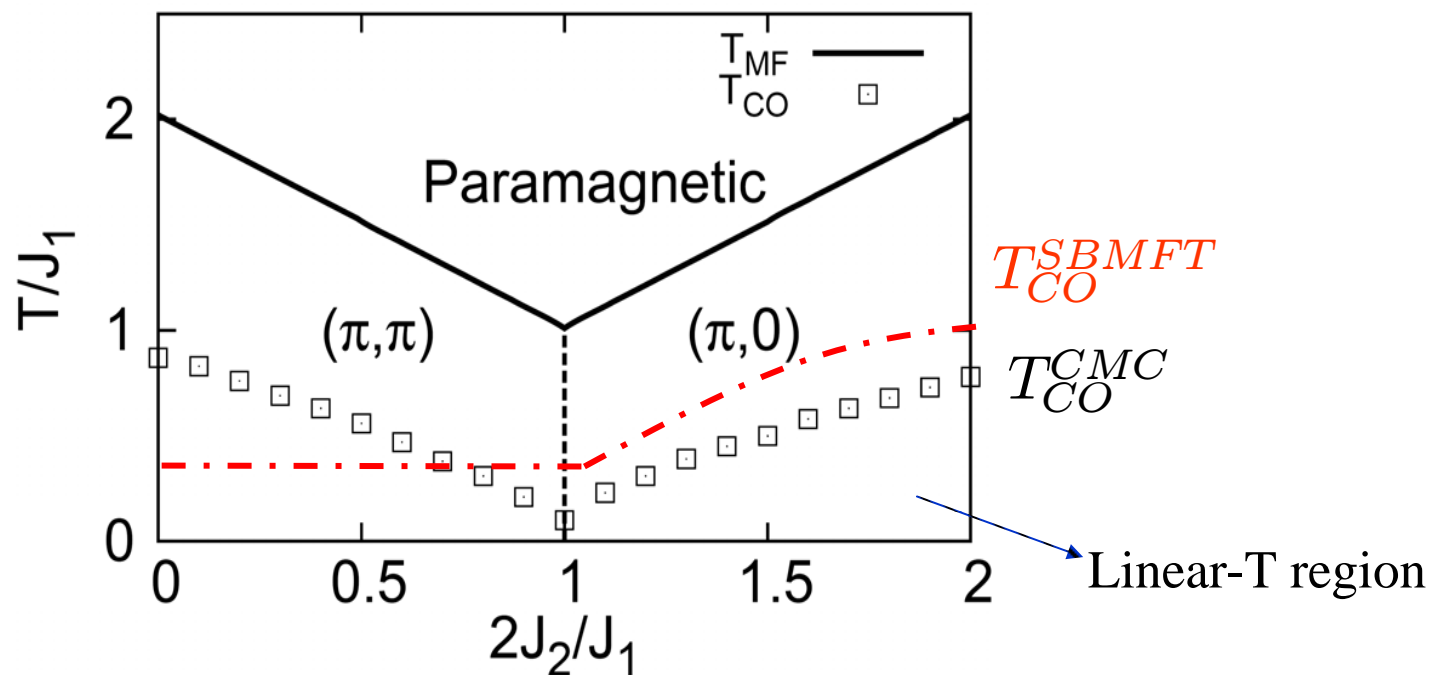
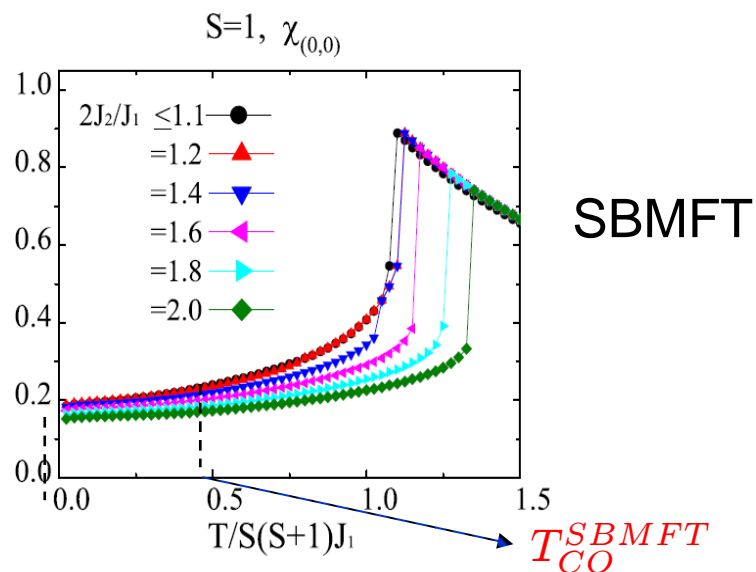
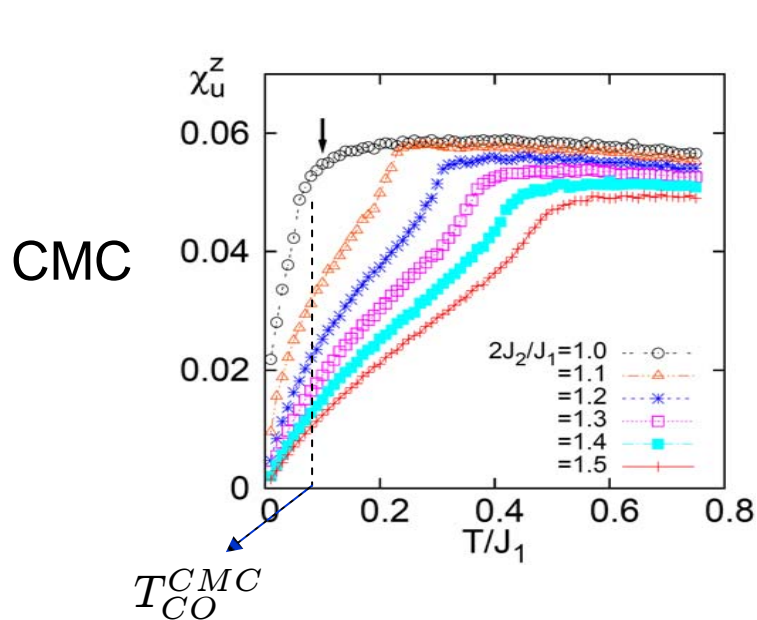
$$T \in [0 \rightarrow (0.3J_1 \sim 0.9J_1)] \quad (S = \frac{1}{2})$$

$$T \in [0 \rightarrow (0.4J_1 \sim 1.0J_1)] \quad (S = 1)$$

- **In the J_1 only limit**

$$A = \frac{2}{C^2}, \quad C \sim 3.2, \quad A \sim 0.19; \quad B = \frac{\ln(L/a)}{\pi \Delta_2^2}$$

CMC vs. SBMFT: uniform susceptibility & phase diagram

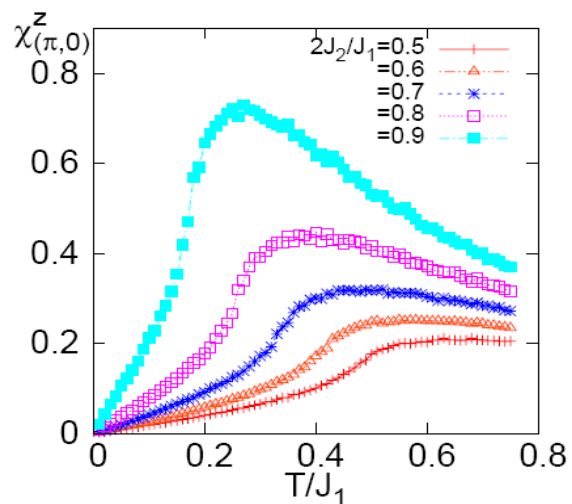


CMC vs. SBMFT: staggered susceptibility

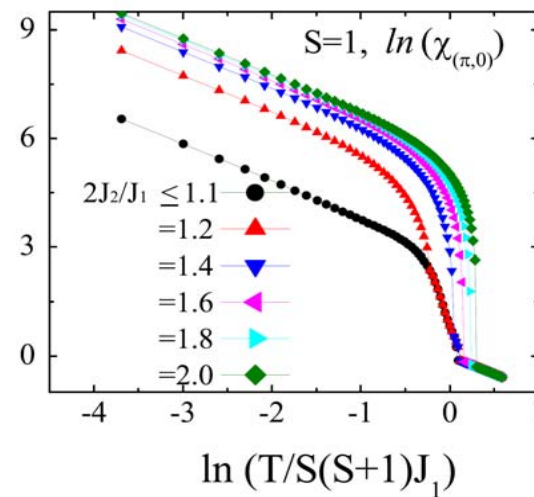
in (π, π) order

$\chi(\pi, 0)$

CMC

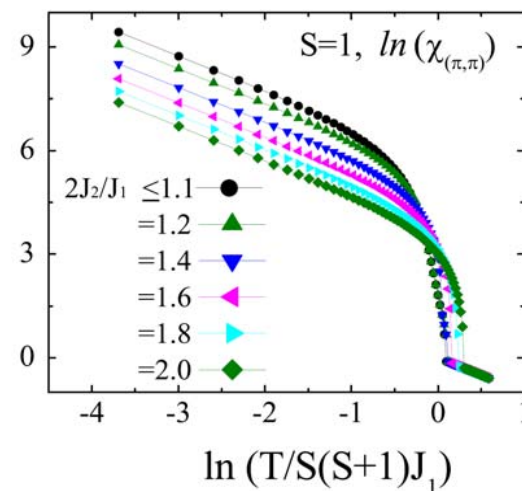
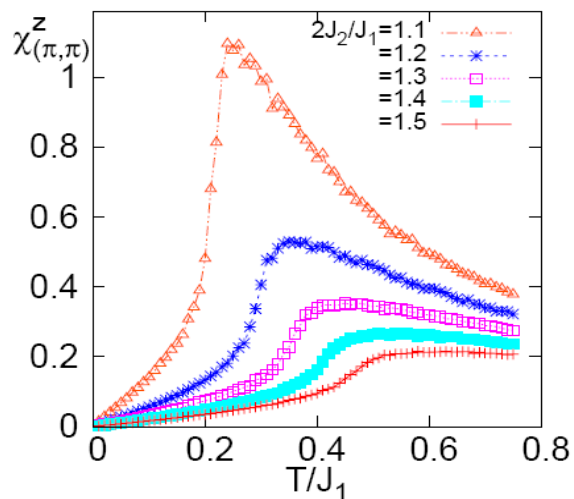


SBMFT



in $(\pi, 0)$ order

$\chi(\pi, \pi)$



Summary

- We calculated the susceptibilities of J_1 - J_2 model in $0 < 2J_2/J_1 < 2$
classical Monte Carlo(CMC) + Schwinger Boson mean field theory(SBMFT)
- The uniform susceptibilities by the two methods are consistent
 - $0 < 2J_2/J_1 < 2$, uniform susceptibility--linear T dependence
 - largely independent of the degree of frustration in the spin model
 - a wider linear temperature range when $2J_2/J_1$ is further removed from the critical point ,i.e. away from the critical point, $T \sim 0.8J_1 - J_1 \rightarrow 400\text{K} \sim 500\text{K}$
 - in agreement with the persistence of linear susceptibility in FeAs compounds.
- Future:local moments & itinerant electrons

Itinerant electron model

